# SPECTRUM OF PARTIAL INTEGRAL OPERATORS WITH DEGENERATE KERNEL

ESHKABILOV YU.KH., ARZIKULOV G.P., AND HAYDAROV F.H.

ABSTRACT. In the paper we consider self-adjoint partial integral operators of Fredholm type T with a degenerate kernel on the space  $L_2([a,b] \times [c,d])$ . Essential and discrete spectra of T are described.

# 1. Introduction

Linear equations and operators involving partial integrals appear in elasticity theory [1, 2, 3], continuum mechanics [2, 4, 5, 6], aerodynamics [7] and in PDE theory [8, 9]. Self-adjoint partial integral operators arise in the theory of Schrödinger operators [10, 11, 12, 13]. Spectrum of a discrete Schrödinger operator H are tightly connected (see [13, 14]) with that of partial integral operators which participate in the presentation of the operator H.

Let  $\Omega_1$  and  $\Omega_2$  be closed boundary subsets in  $\mathbb{R}^{\nu_1}$  and  $\mathbb{R}^{\nu_2}$ , respectively. Partial integral operator (PIO) of Fredholm type in the space  $L_p(\Omega_1 \times \Omega_2)$ ,  $p \geq 1$  is an operator of the form [15]:

$$(1) T = T_0 + T_1 + T_2 + K,$$

where operators  $T_0$ ,  $T_1$ ,  $T_2$  and K are defined by the following formulas:

$$T_{0}f(x,y) = k_{0}(x,y)f(x,y),$$

$$T_{1}f(x,y) = \int_{\Omega_{1}} k_{1}(x,s,y)f(s,y)ds,$$

$$T_{2}f(x,y) = \int_{\Omega_{2}} k_{2}(x,t,y)f(x,t)dt,$$

$$Kf(x,y) = \int_{\Omega_{1}} \int_{\Omega_{2}} k(x,y;s,t)f(s,t)dsdt.$$

Here  $k_0$ ,  $k_1$ ,  $k_2$  and k are given measurable functions on  $\Omega_1 \times \Omega_2$ ,  $\Omega_1^2 \times \Omega_2^2$ ,  $\Omega_1 \times \Omega_2$  and  $(\Omega_1 \times \Omega_2)^2$ , respectively, and all integrals have to be understood in the Lebesgue sense, where  $ds = d\mu_1(s)$ ,  $dt = d\mu_2(t)$ ,  $\mu_k(\cdot)$  is the Lebesgue measure on the  $\sigma$ -algebra of subsets  $\Omega_k$ , k = 1, 2.

In 1975, Likhtarnikov and Vitova [16] spectral properties of partial integral operators are studied. In [16], the following restrictions were imposed:  $k_1(x,s) \in L_2(\Omega_1 \times \Omega_1)$ ,  $k_2(y,t) \in L_2(\Omega_2 \times \Omega_2)$  and  $T_0 = K = 0$ . In [17], spectral properties of PIO with positive kernels were studied (under restriction  $T_0 = K = 0$ ). In Kalitvin and Zabrejko [18] spectral properties of PIO with kernels of two variables

1

 $<sup>2010\</sup> Mathematics\ Subject\ Classification.$  Primary: 45A05, 47A10, 47G10; Secondary: 45P05, 45B05, 45C05.

Key words and phrases. spectrum, essential spectrum, discrete spectrum, partial integral operator, partial integral equation.

in  $L_p, p \ge 1$  are studied. In [19, 20, 21, 22, 23] for the more general equations with continuous kernels or kernels in  $C(L_1)$  spectral properties of the PIO and solvability of partial integral equations in the space  $C([a,b] \times [c,d])$  were studied.

Self-adjoint PIO with  $T_0 \neq 0$  were first studied in [10], where theorem above essential spectrum is proved. Finiteness and infiniteness of a discrete spectrum of self-adjoint PIO, arising in the theory of Schrödinger operators, are investigated in [11, 12, 14]. In [24] applications of partial integral equations and operators to problems of continuous mechanics, elasticity problems and other problems were considered. Still, some important spectral properties of PIO in the space  $L_2$  are left open. The present paper is dedicated to the mentioned problem for PIO with degenerate kernels from the class  $L_2$ .

Let  $T_1$  be a linear integral operator in the space  $L_2([a,b]\times [c,d])$  given by the formula

(2) 
$$(T_1 f)(x,y) = \int_a^b k(x,s,y) f(s,y) ds.$$

Here k(x, s, y) is a measurable function on  $[a, b]^2 \times [c, d]$ .

The kernel k(x, s, y) of the integral operator  $T_1$  usually satisfies the condition

$$\int_{a}^{b} k(x, s, y) f(s, y) ds \in L_{2}([a, b] \times [c, d]), \quad \forall f \in L_{2}([a, b] \times [c, d]).$$

Consequently, the operator  $T_1$  is a linear bounded operator on  $L_2([a,b] \times [c,d])$ . If, in addition, the kernel k(x,s,y) satisfies the condition:

$$k(x, s, y) = \overline{k(s, x, y)},$$
 for almost all  $y \in [c, d],$ 

then the operator  $T_1$  is a self-adjoint operator on the Hilbert space  $L_2([a,b] \times [c,d])$ . Let  $\{\varphi_k(x)\}_{k=1}^n$  be a orthonormal system of functions from the  $L_2[a,b]$ , and let  $\{h_k(y)\}_{k=1}^n$  be a system of essential bounded real functions on [c,d].

We define a measurable function  $k_1(x, s, y)$  on  $[a, b]^2 \times [c, d]$  by the following rule:

(3) 
$$k_1(x,s,y) = \sum_{k=1}^n \varphi_k(x) \overline{\varphi_k(s)} h_k(y).$$

Then the PIO  $T_1$  with the kernel  $k_1(x, s, y)$  is a self-adjoint bounded linear operator on  $L_2([a, b] \times [c, d])$ .

Let  $\{\psi_k(y)\}_{k=\overline{1,m}}$  be a some orthonormal system of functions from the  $L_2[c,d]$ ,  $\{p_k(x)\}_{k=\overline{1,m}}$  be a system of essential bounded real functions on [a,b]. We define the measurable function  $k_2(x,t,y)$  on  $[a,b]\times [c,d]^2$  by the following rule:

(4) 
$$k_2(x,t,y) = \sum_{j=1}^m p_j(x)\psi_j(y)\overline{\psi_j(t)}.$$

Then the PIO  $T_2$  with the kernel  $k_2(x,t,y)$ :

$$T_2 f(x,y) = \int_c^d k_2(x,t,y) f(x,t) d\mu_2(t)$$

is a linear bounded self-adjoint operator on  $L_2([a,b] \times [c,d])$ .

For an essential bounded function  $\varphi \geq 0$  on the measurable set  $\Omega \subset \mathbb{R}^{\nu}$ , we define

$$esssup_{\Omega}(\varphi)=\inf\{C:\mu(\{\xi\in\Omega:\varphi(\xi)>C\})=0\},$$

where  $\mu(\cdot)$  is the Lebesgue measure on  $\mathbb{R}$ . For a measurable function  $\varphi$  on the set  $\Omega \subset \mathbb{R}^{\nu}$ , the number  $\lambda \in \mathbb{R}$  is called an essential value of the function  $\varphi$  if

$$\mu(\{\xi \in \Omega : \lambda - \varepsilon < \varphi(\xi) < \lambda + \varepsilon\}) > 0$$

for all  $\varepsilon > 0$ . We denote by  $Essran(\varphi)$  the set of all essential values of the function  $\varphi$ .

In this paper, we study essential and discrete spectra of PIO of the form  $T_1 + T_2$  with the degenerate kernels. The resolvent set, spectrum, essential spectrum and discrete spectrum are denoted by  $\rho$ ,  $\sigma$ ,  $\sigma_{ess}$  and  $\sigma_{disc}$ ,, respectively (see [25]).

# 2. Spectral property of PIO $T_1$ and $T_2$

In this section, we study spectra of PIO  $T_1$  and  $T_2$ .

**Proposition 2.1.** Zero is an eigenvalue of  $T_1$  of infinite multiplicity. A number  $\lambda_0 \neq 0$  is an eigenvalue of  $T_1$  if and only if there exists  $1 \leq j_0 \leq n$  such that  $\mu_2(h_{j_0}^{-1}(\{\lambda_0\})) > 0$ .

We denote by M the subspace of a Hilbert space  $L_2[a,b]$  constructed by the orthogonal system  $\{\varphi_1,...,\varphi_n\}$ . Then  $\dim M=n$ , and for the subspace  $\mathcal{H}=L_2[a,b]\ominus M$ , we have  $\dim \mathcal{H}=\infty$ .

Let  $\{g_k\}_{k\in\mathbb{N}}$  be an orthonormal basis in  $\mathcal{H}$ , and  $\psi \in L_2[c,d]$ ,  $\|\psi\| = 1$ . Then  $f_k(x,y) = g_k(x)\psi(y) \in L_2([a,b] \times [c,d])$ ,  $k \in \mathbb{N}$ , and the system  $\{f_k\}_{k\in\mathbb{N}}$  is orthonormal in  $L_2([a,b] \times [c,d])$ . Clearly,

$$T_1 f_k = \sum_{i=1}^n \int_a^b \varphi_i(x) \overline{\varphi_i(s)} f_k(s, y) d\mu_1(s) = 0, \ k \in \mathbb{N},$$

i.e. zero is an eigenvalue of PIO  $T_1$  of infinite multiplicity.

"if" part. Let  $\lambda_0 \in \mathbb{C}\setminus\{0\}$  be an eigenvalue of the PIO  $T_1$ . Then there exists  $f \in L_2([a,b] \times [c,d]), ||f|| = 1$  such that

$$T_1 f = \lambda_0 f$$
.

We define compact self-adjoint integral operators  $K_{\omega}$  in  $L_2[a,b]$  as follows:

$$K_{\omega}\varphi(x) = \int_{a}^{b} k_{1}(x, s, w)\varphi(s)d\mu_{1}(s), \ \omega \in \Omega_{0},$$

where

$$\Omega_0 = \{ \omega \in [c, d] : p_{\omega}(x, s) = k_1(x, s, \omega) \in L_2([a, b]^2) \}.$$

We have  $\mu_2([c,d] \setminus \Omega_0) = 0$ . Put

$$\mathcal{M}_1 = \{ \omega \in \Omega_0 : f_{\omega}(x) = f(x, \omega) \in L_2[a, b] \}.$$

Then  $\mu_2(\mathcal{M}_1) > 0$ .

We define measurable subsets:

$$\mathcal{D}_k = \{ \omega \in [c, d] : h_k(\omega) = \lambda_0 \}, \ k \in \{1, ..., n\}.$$

Put

$$\mathcal{D}_0 = \bigcup_{k=1}^n \mathcal{D}_k.$$

Let  $\omega \in \mathcal{M}_1$ . Then  $K_{\omega} f_{\omega} = \lambda_0 f_{\omega}$ , i.e. the number  $\lambda_0$  is an eigenvalue of operators  $K_{\omega}$ ,  $\omega \in \mathcal{M}_1$ . We define by  $\left\{\alpha_1^{(\omega)}, ..., \alpha_{n_{\omega}}^{(\omega)}\right\}$  the set of eigenvalues of the operator  $K_{\omega}$  which are different from zero. Then

$$\lambda_0 \in \{\alpha_1^{(\omega)}, ..., \alpha_{n_{\omega}}^{(\omega)}\} \subset \{h_1(\omega), ..., h_n(\omega)\},\$$

i.e. there exists  $j_0 \in \{1, ..., n\}$  such that  $h_{j_0}(\omega) = \lambda_0$ .

Consequently, we have  $\omega \in \mathcal{D}_0$ . Thus,  $\mathcal{M}_1 \subset \mathcal{D}_0$ . It means that  $\mu_2(\mathcal{D}_0) > 0$ . Then there exists  $j_0 \in \{1, ..., n\}$  such that

$$\mu_2(\{\omega \in [c,d] : h_{j_0}(\omega) = \lambda_0\}) > 0.$$

"only if" part. Let for a  $j_0 \in \{1, ..., n\}$  we have  $\mu_2(h_{j_0}^{-1}(\{\lambda_0\})) > 0$ . Put  $\mathcal{D} = h_{j_0}^{-1}(\{\lambda_0\})$ . We define the function  $\psi \in L_2[c,d]$  in the following way:

$$\psi(y) = \frac{\chi_{\mathcal{D}}(y)}{\sqrt{\mu_2(\mathcal{D})}}, \ y \in [c, d],$$

where  $\chi_G(\cdot)$  is the characteristic function of a set G. Obviously,  $\|\psi\|=1$ . Let  $f(x,y)=\varphi_{j_0}(x)\psi(y)$ . Then  $f\in L_2([a,b]\times[c,d])$  and  $\|f\|=1$ . On the other hand,

$$T_1 f(x,y) = \sum_{k=1}^{n} \varphi_k(x) \int_a^b \overline{\varphi_k(s)} h_k(y) \varphi_{j_0}(s) \psi(y) d\mu_1(s) = \varphi_{j_0}(x) h_{j_0}(y) \psi(y) = \lambda_0 f(x,y),$$

i.e. the number  $\lambda_0$  is an eigenvalue of  $T_1$ .

We consider the following projectors  $P_k$  in the space  $L_2([a,b] \times [c,d])$ :

$$P_k f(x,y) = \int_a^b \varphi_k(x) \overline{\varphi_k(s)} f(s,y) d\mu_1(s), \quad k \in \{1,...,n\}.$$

Let  $P = P_1 + ... + P_n$  and  $P_0 = E - P$ , where E is the identical operator. Then projectors  $P_i$  and  $P_j$   $(i \neq j)$  are orthogonal.

**Proposition 2.2.** If  $\lambda \neq 0$  and  $\lambda \in \bigcup_{k=1}^{n} Essran(h_k)$ , then the operator  $T_1 - \lambda E$  is invertible in  $L_2([a,b] \times [c,d])$ , and the operator  $(T_1 - \lambda E)^{-1}$  is bounded in  $L_2([a,b] \times [c,d])$ , moreover

$$(T_1 - \lambda E)^{-1} f(x, y) = -\frac{1}{\lambda} \left( f(x, y) - \sum_{k=1}^n \frac{h_k(y)}{h_k(y) - \lambda} P_k f(x, y) \right).$$

*Proof.* Let  $\lambda \neq 0$  and  $\lambda \in \bigcup_{k=1}^n Esstan(h_k)$ . We define the operator  $B_{\lambda}$  on the  $L_2([a,b] \times [c,d])$  by the formula:

$$B_{\lambda}f(x,y) = \sum_{k=1}^{n} \frac{1}{h_{k}(y) - \lambda} P_{k}f(x,y) - \frac{1}{\lambda} P_{0}f(x,y).$$

It is clear,

(5) 
$$(T_1 - \lambda E)B_{\lambda} = B_{\lambda}(T_1 - \lambda E) = E.$$

For all  $\lambda \in Essran(h_k) \cup \{0\}$ , the operator

$$A_k f(x,y) = \frac{1}{h_k(y) - \lambda} P_k f(x,y), \ f \in L_2([a,b] \times [c,d])$$

is a bounded operator in  $L_2([a,b] \times [c,d])$ . Then the operator  $B_{\lambda}$  is a bounded operator in  $L_2([a,b] \times [c,d])$ . For each  $\lambda \in \{0\} \cup \left(\bigcup_{k=1}^n Essran(h_k)\right)$ , (5) implies

$$(T - \lambda E)^{-1} = B_{\lambda}.$$

Hence, we have

$$B_{\lambda}f(x,y) = -\frac{1}{\lambda} \left( f(x,y) - \sum_{k=1}^{n} \frac{h_k(y)}{h_k(y) - \lambda} P_k f(x,y) \right).$$

**Theorem 2.3.** For the spectra  $\sigma(T_1)$  of the PIO  $T_1$  with a degenerate kernel (3), the following formula holds:

$$\sigma(T_1) = \{0\} \cup \left(\bigcup_{k=1}^n Essran(h_k)\right).$$

*Proof.* By proposition 2.2, we obtaine

$$\sigma(T_1) \subset \{0\} \cup \left(\bigcup_{k=1}^n Essran(h_k)\right).$$

However, by proposition 2.1 we have  $0 \in \sigma(T_1)$ . Now we prove

$$\bigcup_{k=1}^{n} Essran(h_k) \subset \sigma(T_1).$$

Let  $\lambda_0 \in Essran(h_{j_0}), \ \lambda_0 \neq 0$  and  $t_0$  be arbitrary point from the subset  $h_{j_0}^{-1}(\{\lambda_0\})$ . Put

$$V_i = \left\{ t \in [c, d] : \frac{1}{i+1} < |t_0 - t| < \frac{1}{i} \right\}, \ i \in \mathbb{N}.$$

Then there exists  $n_0 \in \mathbb{N}$  such that  $\mu_2(V_i) > 0$  for all  $i \geq n_0$ . We consider the following sequence of orthonormal functions  $\chi_p(y) \in L_2[c,d]$ :

$$\chi_p(y) = \begin{cases} \frac{1}{\sqrt{\mu_2(V_p)}}, & y \in V_p, \\ 0, & y \in V_p, \end{cases}$$

where  $p \ge n_0$ . We define by  $f_p(x,y) \in L_2([a,b] \times [c,d])$  the orthonormal system of functions:  $f_p(x,y) = \varphi_{j_0}(x)\chi_p(y), \ p \ge n_0$ . Then we have

$$(T_1 - \lambda_0 E) f_p(x, y) = (h_{i_0}(y) - \lambda_0) f_p(x, y).$$

Hence,

$$\|(T_1 - \lambda_0 E)f_p\| \le \sqrt{esssup_{V_p} (h_{j_0}(y) - \lambda_0)^2}, \ p \ge n_0.$$

Since zero is an essential value of the function  $h_{j_0}(y) - \lambda_0$ , then for large  $n_1 \ge n_0$ , there exists a small number  $\delta_{n_1}$  such that

$$|h_{j_0}(y) - \lambda_0| < \delta_{n_1}$$
 for almost all  $y \in V_p$ ,  $p \ge n_1$ .

Therefore

$$||(T_1 - \lambda_0 E) f_n|| < \delta_{n_1}, \quad p > n_1,$$

i.e.  $\lim_{n\to\infty} \| (T_1 - \lambda_0 E) f_n \| = 0$ . This and the Weyl criterion for an essential spectrum of self-adjoint operators [25] imply  $\lambda_0 \in \sigma_{ess}(T_1) \subset \sigma(T_1)$ .

**Proposition 2.4.** Any eigenvalue of the PIO  $T_1$  is of infinite multiplicity.

*Proof.* Let  $\lambda \in \mathbb{R} \setminus \{0\}$  be an eigenvalue of  $T_1$ . Then there exists  $f_0 \in L_2([a,b] \times [c,d])$ ,  $||f_0|| = 1$  such that  $T_1 f_0 = \lambda f_0$ . We define the measurable subset  $\Omega_0 \subset [c,d]$ :

$$\Omega_0 = \left\{ y \in [c, d] : \int_a^b |f_0(x, y)|^2 d\mu_1(x) \neq 0 \right\}.$$

Obviously,  $\mu_2(\Omega_0) > 0$ . Define the function:

$$f_0(x,y) = \begin{cases} \frac{f_0(x,y)}{\sqrt{\int_a^b |f_0(s,y)|^2 d\mu_1(s)}}, & x \in [a,b], y \in \Omega_0 \\ 0, & x \in [a,b], y \in \Omega_0, \end{cases}$$

Then  $f_0 \in L_2([a,b] \times [c,d])$ , and  $f_0 \neq 0$ .

Let  $\{\psi_k\}_k \in \mathbb{N}$  be a system of orthonormal functions from  $L_2(\Omega_0)$ . Consider the sequence of functions from  $L_2([a,b]\times [c,d])$ :

$$f_k(x,y) = f_0(x,y)\psi_k(y), k \in \mathbb{N},$$

where

$$\psi_k(y) = \begin{cases} \widetilde{\psi_k}(y), & y \in \Omega_0 \\ 0, & y \in \Omega_0. \end{cases}$$

Then

$$\int_{a}^{b} \int_{c}^{d} |f_{n}(x,y)|^{2} d\mu_{1}(x) d\mu_{2}(y) = \int_{\Omega_{0}} \left| \widetilde{\psi}_{n}(y) \right|^{2} d\mu_{2}(y) = 1,$$

and

$$(f_i, f_j) = \int_{\Omega_0} \widetilde{\psi}_i(y) \overline{\widetilde{\psi}_j(y)} d\mu_2(y) = 0$$

for  $i \neq j$ . Clearly

$$T_1 f_k(x, y) = \lambda f_k(x, y), \ k \in \mathbb{N},$$

i.e. the number  $\lambda$  is an eigenvalue of the PIO  $T_1$  of infinite multiplicity.

**Corollary 2.5.** A discrete spectrum of the PIO  $T_1$  with a degenerate kernel (3) is absent.

**Corollary 2.6.** If every function  $h_k$ ,  $k \in \{1, ..., n\}$  is continuous and strictly monotone on [c, d], then there is not an eigenvalue of the PIO  $T_1$  different from zero.

Consider the following projectors  $Q_j$  in the space  $L_2([a,b] \times [c,d])$ :

$$Q_{j}f(x,y) = \int_{c}^{d} \psi_{j}(y)\overline{\psi_{j}(t)}f(x,t)d\mu_{2}(t), \quad j \in \{1,...,m\}.$$

**Proposition 2.7.** If  $\lambda \neq 0$  and  $\lambda \in \bigcup_{j=1}^{m} Essran(p_j)$ , then the operator  $T_2 - \lambda E$  is invertible on  $L_2([a,b] \times [c,d])$ , and the operator  $(T_2 - \lambda E)^{-1}$  is bounded in  $L_2([a,b] \times [c,d])$ , moreover

$$(T_2 - \lambda E)^{-1} f(x, y) = -\frac{1}{\lambda} \left( f(x, y) - \sum_{j=1}^m \frac{p_j(x)}{p_j(x) - \lambda} Q_j f(x, y) \right).$$

**Theorem 2.8.** For the spectrum  $\sigma(T_2)$  of the PIO  $T_2$  with a degenerate kernel (4), the following formula is hold:

$$\sigma(T_2) = \{0\} \cup \left(\bigcup_{j=1}^m Essran(p_j)\right).$$

#### 3. Solvability of partial integral equations

We consider the Fredholm partial integral equation (PIE) of the second kind

(6) 
$$f(x,y) - \tau(T_1 + T_2)f(x,y) = g(x,y)$$

in the Hilbert space  $L_2([a,b]\times[c,d])$ , where f is an unknown function from  $L_2([a,b]\times[c,d])$ ,  $g\in L_2([a,b]\times[c,d])$  is a given function, and  $\tau\in\mathbb{C}$  is the parameter of the equation.

The homogeneous PIE corresponding to (6) has the following form:

$$f(x,y) - \tau(T_1 + T_2)f(x,y) = 0.$$

In this section, we reduce to some necessary results for PIE of the second kind. Assume that  $\tau \neq 0$  and  $\tau^{-1} \in \rho(T_1)$ . Then the operator  $E - \tau T_1$  is invertible in  $L_2([a,b] \times [c,d])$ , and the operator  $(E - \tau T_1)^{-1}$  is bounded on  $L_2([a,b] \times [c,d])$ , moreover, by proposition 2.2, we have

$$(E - \tau T_1)^{-1} f(x, y) = f(x, y) + \tau \sum_{k=1}^{n} \frac{h_k(y)}{1 - \tau h_k(y)} P_k f(x, y).$$

Analogously, if  $\tau \neq 0$  and  $\tau^{-1} \in \rho(T_2)$ , then the operator  $E - \tau T_2$  is invertible in  $L_2([a,b] \times [c,d])$ , and the operator  $(E - \tau T_2)^{-1}$  is bounded on  $L_2([a,b] \times [c,d])$ , moreover, by proposition 2.7,

$$(E - \tau T_2)^{-1} f(x, y) = f(x, y) + \tau \sum_{j=1}^{m} \frac{p_j(x)}{1 - \tau p_j(x)} Q_j f(x, y).$$

Let  $\tau \neq 0$  and  $\tau^{-1} \in \rho(T_1) \cap \rho(T_2)$ . We define compact operators  $W_1(\tau)$  and  $W_2(\tau)$  by the following formulas:

$$W_1(\tau) = (E - \tau T_2)^{-1} S_1(\tau) T_2, \quad W_2(\tau) = (E - \tau T_1)^{-1} S_2(\tau) T_1,$$

where

$$S_1(\tau)f(x,y) = \sum_{k=1}^n \frac{h_k(y)}{1 - \tau h_k(y)} (P_k f)(x,y), \quad S_2(\tau)f(x,y) = \sum_{i=1}^m \frac{p_i(x)}{1 - \tau p_i(x)} (Q_i f)(x,y).$$

**Lemma 3.1.** Let  $\tau \neq 0$  and  $\tau^{-1} \in \rho(T_1) \cap \rho(T_2)$ . Then the following three homogenous Fredholm PIE of the second kind are equivalent:

(7) 
$$f - \tau (T_1 + T_2)f = 0,$$

(8) 
$$f - \tau^2 W_1(\tau) f = 0,$$

(9) 
$$f - \tau^2 W_2(\tau) f = 0.$$

In [19, 20, 21, 22, 23] lemmas similar to the lemma 3.1 for the case of general PIE's in  $C([a,b] \times [c,d])$  with continuous kernels or kernels in  $C(L_1)$  were proved. The scheme for the proof of lemma 3.1 can be seen from these works.

Assume that  $\tau \neq 0$  and  $\tau^{-1} \in \rho(T_1) \cap \rho(T_2)$ . We denote by  $\Delta_1(\tau)$  and  $\Delta_2(\tau)$  the Fredholm determinants of the operators  $E - \tau^2 W_1(\tau)$  and  $E - \tau^2 W_2(\tau)$ , respectively. Define in  $\mathbb C$  the following subsets

$$\mathcal{R}_1 = \left\{ \tau \in \mathbb{C} \setminus \{0\} : \tau^{-1} \in \rho(T_1) \cap \rho(T_2) \text{ and } \Delta_1(\tau) \neq 0 \right\},$$

$$\mathcal{R}_2 = \left\{ \tau \in \mathbb{C} \setminus \{0\} : \tau^{-1} \in \rho(T_1) \cap \rho(T_2) \text{ and } \Delta_2(\tau) \neq 0 \right\},$$

$$\mathcal{D}_1 = \left\{ \tau \in \mathbb{C} \setminus \{0\} : \tau^{-1} \in \rho(T_1) \cap \rho(T_2) \text{ and } \Delta_1(\tau) = 0 \right\},$$

$$\mathcal{D}_2 = \left\{ \tau \in \mathbb{C} \setminus \{0\} : \tau^{-1} \in \rho(T_1) \cap \rho(T_2) \text{ and } \Delta_2(\tau) = 0 \right\}.$$

It follows from lemma 3.1 that  $\mathcal{R}_1 = \mathcal{R}_2$  and  $\mathcal{D}_1 = \mathcal{D}_2$ . Put

$$\mathcal{R} = \mathcal{R}(T) = \mathcal{R}_1$$
 and  $\mathcal{D} = \mathcal{D}(T) = \mathcal{D}_1$ .

Then we obtain

(10) 
$$\mathcal{R} \cup \mathcal{D} = \left\{ \tau \in \mathbb{C} : \tau \in \mathbb{C} \setminus \{0\} \text{ and } \tau^{-1} \in \rho(T_1) \cap \rho(T_2) \right\}$$

and  $\mathcal{R} \cap \mathcal{D} = \emptyset$ .

Lemma 3.1 implies the following

**Theorem 3.2.** Let  $\tau \in \mathcal{R} \cup \mathcal{D}$ . Homogenous PIE (7) has a non-trivial solution if and only if  $\tau \in \mathcal{D}$ .

**Theorem 3.3.** Let  $\tau \in \mathcal{R}$ . Then the Fredholm PIE of the second kind

(11) 
$$f - \tau (T_1 + T_2)f = g$$

has the unique solution  $f_0 \in L_2([a,b] \times [c,d])$  for any  $g \in L_2([a,b] \times [c,d])$ .

4. Spectrum of the PIO  $T_1 + T_2$ 

We put

$$\mathcal{D}_0 = \mathcal{D}_0(T) = \left\{ \xi : \xi = \frac{1}{\tau}, \ \tau \in \mathcal{D}(T) \right\}.$$

**Lemma 4.1.** For each  $\lambda \in \mathbb{C} \setminus (\sigma(T_1) \cup \sigma(T_2) \cup \mathcal{D}_0(T))$ , the resolvent  $R_{\lambda}(T)$  of the PIO  $T = T_1 + T_2$  exists and is bounded on  $L_2([a, b] \times [c, d])$ .

*Proof.* Let  $\lambda \in \mathbb{C} \setminus (\sigma(T_1) \cup \sigma(T_2) \cup \mathcal{D}_0(T))$ . Then  $\lambda \neq 0$ , and  $\lambda \in \rho(T_1) \cap \rho(T_2)$ . Then operators  $E - \frac{1}{\lambda}T_1$  and  $E - \frac{1}{\lambda}T_2$  are injective and

$$\left(E - \frac{1}{\lambda}T_1\right)^{-1} = E + \frac{1}{\lambda}S_1\left(\frac{1}{\lambda}\right), \quad \left(E - \frac{1}{\lambda}T_2\right)^{-1} = E + \frac{1}{\lambda}S_2\left(\frac{1}{\lambda}\right).$$

However, from  $\lambda \in \mathcal{D}_0(T)$  and by (10), we obtain  $\Delta_1\left(\frac{1}{\lambda}\right) \neq 0$ . Consequently, the operator  $E - \frac{1}{\lambda^2}W_1\left(\frac{1}{\lambda}\right)$  is injective. By the other hand, we have

$$T_1 + T_2 - \lambda E = (T_1 - \lambda E)(E + (T_1 - \lambda E)^{-1}T_2) =$$

$$= (T_1 - \lambda E) \left[ E - \frac{1}{\lambda} \left( E + \frac{1}{\lambda} S_1 \left( \frac{1}{\lambda} \right) \right) T_2 \right] =$$

$$= (T_1 - \lambda E) \left( E - \frac{1}{\lambda} T_2 \right) \left( E - \frac{1}{\lambda^2} \left( E - \frac{1}{\lambda} T_2 \right)^{-1} S_1 \left( \frac{1}{\lambda} \right) T_2 \right) =$$

$$= -\lambda \left( E - \frac{1}{\lambda} T_1 \right) \left( E - \frac{1}{\lambda} T_2 \right) \left( E - \frac{1}{\lambda^2} W_1 \left( \frac{1}{\lambda} \right) \right).$$

Hence,

$$R_{\lambda}(T) = -\frac{1}{\lambda} \left( E - \frac{1}{\lambda^2} W_1 \left( \frac{1}{\lambda} \right) \right)^{-1} \left( E - \frac{1}{\lambda} T_2 \right)^{-1} \left( E - \frac{1}{\lambda} T_1 \right)^{-1}.$$

Boundedness of the operator  $R_{\lambda}(T)$  follows from the last equality.

Corollary 4.2. For the resolvent operator  $R_{\lambda}(T)$  of the PIO  $T=T_1+T_2$ , the formula

$$R_{\lambda}(T) = -\frac{1}{\lambda} \left( E - \frac{1}{\lambda^2} W_1 \left( \frac{1}{\lambda} \right) \right)^{-1} \left( E - \frac{1}{\lambda} T_2 \right)^{-1} \left( E - \frac{1}{\lambda} T_1 \right)^{-1}$$

holds for each  $\lambda \in \mathbb{C} \setminus (\sigma(T_1) \cup \sigma(T_2) \cup \mathcal{D}_0(T))$ .

**Proposition 4.3.** Zero is an eigenvalue of infinite multiplicity of the PIO  $T = T_1 + T_2$ .

*Proof.* Define by  $\mathcal{L}$  the subspace of a Hilbert space  $L_2([a,b] \times [c,d])$  constructed by the orthogonal system  $\{\varphi_i(x)\psi_j(y)\}_{i=\overline{1,n},\ j=\overline{1,m}}$ . Then  $\dim \mathcal{L}=m\times n$  and for subspace  $\mathcal{H}_0=L_2([a,b]\times [c,d])\ominus \mathcal{L}$ , we have  $\dim \mathcal{H}_0=\infty$ .

Let  $\{f_k\}_{k\in\mathbb{N}}$  be an orthonormal basis in  $\mathcal{H}_0$ . It is obvious that

$$T_1 f_p(x,y) = \sum_{i=1}^n \int_a^b \varphi_i(x) \overline{\varphi_i(s)} h_i(y) f_k(s,y) d\mu_1(s) = 0, \ p \in \mathbb{N},$$

$$T_2 f_p(x,y) = \sum_{j=1}^m \int_c^d p_j(x) \psi_j(y) \overline{\psi_j(t)} f_k(x,t) d\mu_2(t) = 0, \ \ p \in \mathbb{N}.$$

i.e.

$$(T_1 + T_2)f_k(x, y) = 0, k \in \mathbb{N}.$$

**Proposition 4.4.** The inclusion  $\sigma(T_1) \cup \sigma(T_2) \subset \sigma_{ess}(T_1 + T_2)$  holds.

*Proof.* We show that  $\sigma(T_1) \subset \sigma_{ess}(T_1 + T_2)$  (the inclusion  $\sigma(T_2) \subset \sigma_{ess}(T_1 + T_2)$  can be proved analogously). We have

$$\sigma(T_1) = \{0\} \cup \left(\bigcup_{i=1}^n Essran(h_i)\right).$$

By proposition 4.3,  $0 \in \sigma_{ess}(T_1 + T_2)$ . Assume that  $\lambda_0 \in Esstan(h_{i_0}), \lambda_0 \neq 0$ , where  $i_0 \in \{1, 2, ..., n\}$  and  $y_0$  is arbitrary point of the set  $h_{i_0}^{-1}(\{\lambda_0\})$ . Put

$$V_k = \left\{ y \in [c, d] : \frac{1}{k+1} < |y_0 - y| < \frac{1}{k} \right\}, \ k \in \mathbb{N}.$$

Then there exists  $k_0 \in \mathbb{N}$  such that  $\mu_2(V_k) > 0$  for all  $k \geq k_0$ , and  $\lim_{n \to \infty} \mu_2(V_k) = 0$ . Consider the sequence of orthonormal functions  $\Phi_k(y) \in L_2[c,d]$ :

$$\Phi_k(y) = \frac{\chi_{V_k}(y)}{\sqrt{\mu_2(V_k)}}, \quad k \in \mathbb{N}, \quad k \ge k_0.$$

Put

$$f_k(x,y) = \varphi_{i_0}(x)\Phi_k(y), \quad k \ge k_0.$$

Then the system of functions  $\{f_k(x,y)\}_{k\geq k_0}$  of  $L_2([a,b]\times [c,d])$  is an orthonormal system.

Now we prove that  $\lim_{k\to\infty} \|(T_1+T_2-\lambda_0 E)f_k\| = 0$ . However,  $\lim_{k\to\infty} \|(T_1-\lambda_0 E)f_k\| = 0$  (see the proof of theorem 2.3). We show that  $\lim_{k\to\infty} \|T_2f_k\| = 0$ .

We define operators  $A_i$ , i = 1, ..., m in the following way:

$$A_i f(x,y) = \int_0^d p_i(x) \psi_i(y) \overline{\psi_i(t)} f(x,t) d\mu_2(t), \quad f \in L_2([a,b] \times [c,d]).$$

Then

$$||A_{i}f_{k}||^{2} \leq \int_{a}^{b} \int_{c}^{d} \left( \int_{c}^{d} |p_{i}(x)| \cdot |\psi_{i}(t)| \cdot |\psi_{i}(y)| \cdot |\varphi_{i_{0}}(x)| \cdot |\Phi_{k}(t)| d\mu_{2}(t) \right)^{2} d\mu_{1}(x) \cdot d\mu_{2}(y) \leq C_{i}^{2} \left( \int_{c}^{d} |\psi_{i}(t)| \cdot \Phi_{k}(t) d\mu_{2}(t) \right)^{2} = \frac{C_{i}^{2}}{\mu_{2}(V_{k})} \left( \int_{V_{k}} |\psi_{i}(t)| d\mu_{2}(t) \right)^{2} \leq C_{i}^{2} \cdot \int_{C} |\psi_{i}(t)|^{2} d\mu(t), \quad i \in \{1, ..., n\},$$

where  $C_i = esssup_{[a,b]} \mid p_i(x) \mid$  . Since Lebesgue integrals are absolute continuous, we obtain

$$\lim_{k \to \infty} \int_{V_k} |\psi_i(t)|^2 d\mu_2(t) = 0$$

at  $\lim_{k\to\infty} \mu_2(V_k) = 0$  and  $\|\psi_i\| = 1$ .

Thus, we get  $\lim_{k\to\infty} ||A_i f_k|| = 0$ ,  $i \in \{1, ..., m\}$ .

However,

$$||T_2 f_k|| \le \sum_{i=1}^m ||A_i f_k||,$$

what follows  $\lim_{k\to\infty} ||T_2f_k|| = 0$ . Hence,  $\lim_{k\to\infty} ||(T_1 + T_2 - \lambda_0 E)f_k|| = 0$ . Finally, by the Weyl criterion for an essential spectrum of self-adjoint operators [25],  $\lambda_0 \in \sigma_{ess}(T_1 + T_2)$ .

**Proposition 4.5.** Each  $\lambda \in \mathcal{D}_0(T)$  is an eigenvalue of finite multiplicity of the PIO  $T = T_1 + T_2$ .

*Proof.* Let  $\lambda \in \mathcal{D}_0(T)$ . Then  $\lambda \neq 0$  and  $\Delta_1\left(\frac{1}{\lambda}\right) = 0$ , where  $\Delta_1(\tau)$  is the Fredholm determinant of the operator  $E - \tau^2 W_1(\tau)$ . It means that the number 1 is an eigenvalue of the compact integral operator  $\frac{1}{\lambda^2} W_1\left(\frac{1}{\lambda}\right)$ . By lemma 3.1, the number  $\lambda$  is an eigenvalue of the PIO  $T_1 + T_2$ . Since the following integral equations

$$f - \frac{1}{\lambda}(T_1 + T_2)f = 0$$

and

$$f - \frac{1}{\lambda^2} W_1 \left(\frac{1}{\lambda}\right) f = 0$$

are equivalent, the number  $\lambda$  is an eigenvalue of finite multiplicity of  $T_1+T_2$  because of every eigenvalue  $\alpha \neq 0$  of compact operators is of finite multiplicity.

The next theorem follows from lemma 4.1 and propositions 4.4, 4.5.

**Theorem 4.6.** For the spectrum  $\sigma(T)$  of the PIO  $T = T_1 + T_2$  with a degenerate kernels, the following formula

$$\sigma(T_1 + T_2) = \{0\} \cup \left(\bigcup_{i=1}^n Essran(h_i)\right) \cup \left(\bigcup_{j=1}^m Essran(p_j)\right) \cup \mathcal{D}_0(T)$$

holds.

# 5. Discrete spectrum of the PIO $T_1 + T_2$

Put  $G = \mathbb{C} \setminus (\sigma(T_1) \cup \sigma(T_2))$ . It is well-known that spectra of a linear bounded self-adjoint operators are compact set in the set of all real numbers. Consequently, the set  $\sigma(T_1) \cup \sigma(T_2)$  is a compact subset in  $\mathbb{R}$ . Therefore the set G is an open subset in  $\mathbb{C}$  and G is unbounded domain in  $\mathbb{C}$ .

For each  $\lambda \in G$ , we consider the kernel of the compact integral operator  $W_1\left(\frac{1}{\lambda}\right)$  given as follows:

$$W_1\left(\frac{1}{\lambda}\right) = \left(E - \frac{1}{\lambda}T_2\right)^{-1} S_1\left(\frac{1}{\lambda}\right)T_2.$$

**Proposition 5.1.** For the kernel  $K(x, y; s, t | \lambda)$  ( $\lambda \in G$ ) of the Fredholm integral operator  $W_1\left(\frac{1}{\lambda}\right)$ , the equality

(12) 
$$\mathcal{K}(x,y;s,t|\lambda) = \lambda \sum_{i=1}^{n} \sum_{k=1}^{m} F_{k,j}(x,y;\lambda) B_{k,j}(s,t),$$

is valid, where (13)

$$F_{k,j}(x,y;\lambda) = \varphi_j(x) \left( \frac{\psi_k(y) h_j(y)}{\lambda - h_j(y)} + \sum_{i=1}^m \frac{p_i(x) \psi_i(y)}{\lambda - p_i(x)} \int_0^d \frac{h_j(\xi)}{\lambda - h_j(\xi)} \psi_k(\xi) \overline{\psi_i(\xi)} d\mu_2(\xi) \right),$$

$$B_{k,j}(s,t) = p_k(s)\overline{\varphi_j(s)} \cdot \overline{\psi_k(t)}.$$

*Proof.* Let  $\lambda \in G$ . Then  $\lambda \in \rho(T_1) \cap \rho(T_2)$ , and we get

$$\left(E - \frac{1}{\lambda}T_2\right)^{-1} = E + \frac{1}{\lambda}S_2\left(\frac{1}{\lambda}\right).$$

However,

$$W_1\left(\frac{1}{\lambda}\right) = S_1\left(\frac{1}{\lambda}\right)T_2 + \frac{1}{\lambda}S_2\left(\frac{1}{\lambda}\right)S_1\left(\frac{1}{\lambda}\right)T_2.$$

For each  $f \in L_2([a, b] \times [c, d])$ , using representations of operators  $S_1(\tau)$  and  $S_2(\tau)$ , we obtain

$$S_1\left(\frac{1}{\lambda}\right)T_2f(x,y) = \sum_{j=1}^n \sum_{k=1}^m \int_a^b \int_a^d K_{k,j}(x,y;\lambda)B_{k,j}(s,t)f(s,t)d\mu_1(s)d\mu_2(t),$$

$$S_{2}\left(\frac{1}{\lambda}\right)S_{1}\left(\frac{1}{\lambda}\right)T_{2}f(x,y) = \sum_{j=1}^{n}\sum_{k=1}^{m}\int_{a}^{b}\int_{a}^{d}G_{k,j}(x,y;\lambda)B_{k,j}(s,t)f(s,t)d\mu_{1}(s)d\mu_{2}(t),$$

where

$$K_{k,j}(x,y;\lambda) = \frac{\lambda \varphi_j(x) \psi_k(y) h_j(y)}{\lambda - h_j(y)},$$

$$G_{k,j}(x,y;\lambda) = \lambda^2 \sum_{i=1}^m \frac{p_i(x)\psi_i(y)}{\lambda - p_i(x)} \int_c^d \frac{h_j(\xi)}{\lambda - h_j(\xi)} \psi_k(\xi) \overline{\psi_i(\xi)} d\mu_2(\xi).$$

Hence, we obtain equality (12) for the kernel  $\mathcal{K}(x, y; s, t | \lambda)$  of the integral operator  $W_1\left(\frac{1}{\lambda}\right)$ .

Set

$$\Gamma_1 = \{1, 2, ..., m\}, \quad \Gamma_2 = \{1, 2, ..., n\} \text{ and } \Gamma = \Gamma_1 \times \Gamma_2.$$

We can introduce the relation of partial order in the set  $\Gamma$  by the following way: for elements  $\omega = (k_1, j_1) \in \Gamma$  and  $\omega' = (k_2, j_2) \in \Gamma$ , we write  $\omega \leq \omega'$  if  $k_1 < k_2$  or  $k_1 = k_2, j_1 \leq j_2$ . As the set  $\Gamma$  is finite, the set  $\Gamma$  is linear complete ordered, i.e. for arbitrary  $\omega, \omega' \in \Gamma$ , we have  $\omega \leq \omega'$  or  $\omega' \leq \omega$ . Thus, we can give elements of  $\Gamma$  in the increase order:

$$\Gamma = \{\omega_1, \omega_2, ..., \omega_{m \cdot (n-1)}, \omega_{m \cdot n}\},\$$

moreover

$$\omega_1 = (1,1) < \omega_2 = (1,2) < \dots < \omega_n = (1,n) < \omega_{n+1} = (2,1) < \dots < \omega_{m:n} = (m,n).$$

Let  $\lambda \in G$  be fixed. For every  $\omega = (k, j) \in \Gamma$ , we define the function  $F_{\omega}(x, y; \lambda)$  on  $[a, b] \times [c, d]$  by the following formula:

$$F_{\omega}(x,y;\lambda) = F_{k,i}(x,y;\lambda).$$

Consider the homogenous Fredholm integral equation

(14) 
$$f(x,y) - \frac{1}{\lambda^2} W_1\left(\frac{1}{\lambda}\right) f(x,y) = 0, \ f \in L_2([a,b] \times [c,d]).$$

Set

$$\int\limits_a^b\int\limits_c^d F_{\omega_i}(x,y;\lambda)f(x,y)d\mu(x)d\mu(y)=A_{\omega_i}(\lambda),\ i\in\{1,...,m\cdot n\}.$$

Then the homogenous equation (14) turns into the equation

$$f(x,y) = \frac{1}{\lambda} \sum_{i=1}^{m \cdot n} A_{\omega_i}(\lambda) F_{\omega_i}(x,y;\lambda).$$

Let

$$\int_{a}^{b} \int_{c}^{d} F_{\omega_{i}}(x, y; \lambda) B_{\omega_{l}}(x, y) d\mu_{1}(x) d\mu_{2}(y) = \Pi_{i, l}(\lambda), \quad i, l \in \{1, ..., m \cdot n\},$$

where

$$B_{\omega_l}(x,y) = B_{k_l,j_l}(x,y), \quad \omega_l = (k_l,j_l).$$

Then we obtain a system of homogenous linear algebraic equations for unknown numbers  $A_{\omega_i}(\lambda)$ :

(15) 
$$A_{\omega_i}(\lambda) - \frac{1}{\lambda} \sum_{l=1}^{m \cdot n} \Pi_{i,l}(\lambda) A_{\omega_l}(\lambda) = 0, \ i \in \{1, ..., m \cdot n\}.$$

**Lemma 5.2.** Let  $\lambda \in G$ . The homogenous integral equation (14) has a nontrivial solution if and only if  $\Delta(\lambda) = 0$ , where

$$\Delta(\lambda) = \begin{vmatrix} \Pi_{1,1}(\lambda) - \lambda & \Pi_{1,2}(\lambda) & \dots & \Pi_{1,m \cdot n}(\lambda) \\ \Pi_{2,1}(\lambda) & \Pi_{2,2}(\lambda) - \lambda & \dots & \Pi_{2,m \cdot n}(\lambda) \\ \dots & \dots & \dots & \dots \\ \Pi_{m \cdot n,1}(\lambda) & \Pi_{m \cdot n,2}(\lambda) & \dots & \Pi_{m \cdot n,m \cdot n}(\lambda) - \lambda \end{vmatrix}.$$

*Proof.* Let  $\lambda \in G$ . Then equivalence of the Fredholm integral equation of the second kind (14) and the system of linear algebraic homogeneous equations (15) is clear. The determinant  $\widetilde{\Delta}(\lambda)$  of the system of equations (15) has the following form:

$$\widetilde{\Delta}(\lambda) = \begin{vmatrix} 1 - \frac{\Pi_{1,1}(\lambda)}{\lambda} & -\frac{\Pi_{1,2}(\lambda)}{\lambda} & \dots & -\frac{\Pi_{1,m\cdot n}(\lambda)}{\lambda} \\ -\frac{\Pi_{2,1}(\lambda)}{\lambda} & 1 - \frac{\Pi_{2,2}(\lambda)}{\lambda} & \dots & -\frac{\Pi_{2,m\cdot n}(\lambda)}{\lambda} \\ \dots & \dots & \dots & \dots \\ -\frac{\Pi_{m\cdot n,1}(\lambda)}{\lambda} & -\frac{\Pi_{m\cdot n,2}(\lambda)}{\lambda} & \dots & 1 - \frac{\Pi_{m\cdot n,m\cdot n}(\lambda)}{\lambda} \end{vmatrix},$$

and

$$\widetilde{\Delta}(\lambda) = \left(-\frac{1}{\lambda}\right)^{m \cdot n} \Delta(\lambda).$$

It is well-known, the system of linear homogeneous equations (15) has nontrivial solution if and only if  $\widetilde{\Delta}(\lambda) = 0$ , i.e.  $\Delta(\lambda) = 0$ . However, we obtain that the homogeneous Fredholm equation (14) has nontrivial solution if and only if  $\Delta(\lambda) = 0$ .

**Lemma 5.3.** The function  $\Delta(z)$  is holomorphic in the domain G.

*Proof.* Let  $\omega \in \Gamma$ . It is known, the function  $F_{\omega}(\lambda) = F_{\omega}(x, y; \lambda)$  is holomorphic by  $\lambda$  in the domain G for almost all  $(x, y) \in [a, b] \times [c, d]$ , and for every  $\lambda \in G$  the integral

$$\int_{a}^{b} \int_{c}^{d} F_{\omega}(x, y; \lambda) B_{\omega'}(x, y) d\mu_{1}(x) d\mu_{2}(y), \quad \omega, \omega' \in \Gamma$$

exists and is finite. Then for every  $\omega = (i, l) \in \Gamma$ , the function  $\Pi_{i,l}(z)$  is a holomorphic function in G. Consequently, the function  $\Delta(z)$  is a sum of holomorphic functions  $F_{\omega_1}(z), F_{\omega_2}(z), ..., F_{\omega_{m \cdot n}}(z)$ , i.e.  $\Delta(z)$  is holomorphic in G.

Remark 5.4. An analogue of Lemma 5.3 can be proved for the general PIE.

**Theorem 5.5.** The discrete spectrum of the PIO  $T = T_1 + T_2$  coincides with the set  $\mathcal{D}_0(T)$ .

*Proof.* Lemmas 3.1 and 5.2 imply

$$\mathcal{D}_0(T) = \{ \lambda \in G : \Delta(\lambda) = 0 \}.$$

By proposition 4.4, we have

$$\sigma(T_1) \cup \sigma(T_2) \subset \sigma_{ess}(T_1 + T_2).$$

By theorem 4.6, we obtain

$$\sigma_{disc}(T) \subset \mathcal{D}_0(T)$$
.

Let  $\lambda_0 \in \mathcal{D}_0(T)$  be arbitrary. Then by proposition 4.5, the number  $\lambda_0$  is an eigenvalue of finite multiplicity of the operator  $T_1 + T_2$ . Since the function  $\Delta(z)$  is holomorphic in the G, arbitrary point  $\lambda$  form the  $\mathcal{D}_0(T)$  is isolate in  $\mathcal{D}_0(T)$ . Then the point  $\lambda_0$  is isolate in the spectrum  $\sigma(T_1) \cup \sigma(T_2) \cup \mathcal{D}_0(T)$  of the operator T since  $(\sigma(T_1) \cup \sigma(T_2)) \cap \mathcal{D}_0(T) = \emptyset$ . Thus, by definition of a discrete spectrum, we obtain  $\lambda_0 \in \sigma_{disc}(T)$ , i.e.  $\mathcal{D}_0(T) \subset \sigma_{disc}(T)$ .

Theorems 4.6 and 5.5 implies

## Theorem 5.6.

$$\sigma_{ess}(T_1 + T_2) = \{0\} \cup \left(\bigcup_{i=1}^n Essran(h_i)\right) \cup \left(\bigcup_{j=1}^m Essran(p_j)\right).$$

**Example 5.7.** Let  $h_i(y) \equiv a_i \in \mathbb{R} \setminus \{0\}$ ,  $i \in \{1, ..., n\}$  and  $p_j(x) \equiv b_j \in \mathbb{R} \setminus \{0\}$ ,  $j \in \{1, ..., m\}$  for the kernels  $k_1(x, s, y)$  and  $k_2(x, t, y)$  of PIO  $T_1$  and  $T_2$ .

Then by theorem 2.3 and 2.8, we obtain

$$\sigma(T_1)=\{0,a_1,...,a_n\}, \quad \sigma(T_2)=\{0,b_1,...,b_m\}.$$

By theorem 5.6, we have

$$\sigma_{ess}(T_1 + T_2) = \{0, a_1, ..., a_n, b_1, ..., b_m\}.$$

We obtain from (13):

$$F_{k,j}(x,y;\lambda) = \frac{\lambda a_j}{(a_j - \lambda)(b_k - \lambda)} \varphi_j(x) \psi_k(y), \quad \lambda \in \sigma_{ess}(T_1 + T_2),$$

and

$$B_{k,j}(s,t) = b_k \overline{\varphi_j(s)} \cdot \overline{\psi_k(t)}.$$

Then the homogeneous Fredholm integral equation (14) becomes the following form (16)

$$f(x,y) - \sum_{j=1}^{n} \sum_{k=1}^{m} \frac{a_j b_k}{(a_j - \lambda)(b_k - \lambda)} \int_a^b \int_c^d \varphi_j(x) \psi_k(y) \overline{\varphi_j(s)} \cdot \overline{\psi_k(t)} d\mu_1(s) d\mu_2(t) = 0.$$

However, using the property of Fredholm integral equations with a degenerate kernel, we obtain

$$\frac{a_j b_k}{(a_j - \lambda)(b_k - \lambda)} = 1, \ j \in \{1, ..., n\}, \ k \in \{1, ..., m\}, \ \lambda \overline{\in} \sigma_{ess}(T_1 + T_2).$$

It means that

$$a_j b_k = (a_j - \lambda)(b_k - \lambda), \ \lambda \overline{\in} \{0\} \cup \{a_i\} \cup \{b_l\},$$

i.e. the integral equation (16) has nontrivial solution if and only if

$$\lambda = a_i + b_k \overline{\in} \{0\} \cup \{a_i\} \cup \{b_l\}.$$

Set

$$\Lambda = \{\lambda : \lambda = a_j + b_k \overline{\in} \{0, a_1, ..., a_n, b_1, ..., b_m\}, \ j = \overline{1, n}, \ k = \overline{1, m}\}.$$

Therefore according to theorem 5.5, we obtain  $\sigma_{disc}(T_1 + T_2) = \Lambda$  and theorem 4.6 implies

(17) 
$$\sigma(T_1 + T_2) = \{\lambda : \lambda = a + b, \ a \in \sigma(T_1), \ b \in \sigma(T_2)\}.$$

It should be noted, the equality (17) was proved for the PIO T in  $L_p, p \ge 1$  with more general kernels  $k_1(x, s, y) = k_1(x, s), k_2(x, t, y) = k_2(t, y)$  in the paper [18], [24].

**Remark 5.8.** It is known, that the discrete spectrum  $\sigma_{disc}(K)$  for each self-adjoint Fredholm integral operator K with a degenerate kernel is finite (since  $\sigma_{disc}(K)$  is the set of all eigenvalues of K different from zero). The following question is arisen: Does this property hold for the PIO  $T = T_1 + T_2$  with a degenerate kernels (3) and (4) This question is still an open problem.

### References

- [1] Kalitvin A.S. On partial integral operators in contact problems of elasticity. (in Russian) Proc. 26 Voronezh Winter School, 1994, 54.
- [2] Kovalenko E.V. On the approximate solution of one type of integral equations arising in elasticity type mathematical physics. (in Russian). Izv. Akad. Nauk Arm. SSR **34**, 5 (1981), pp.14-26.
- [3] Vekua I.N. New Methods of Solving Elliptic Equations. (in Russian) Moscow-Leningrad, Gostekhizdat, 1948.
- [4] Aleksandrov V.M., Kovalenko E.V. On some class of integral equations arising in mixed boundary value problems of continuum mechanics. Sov. Phys. Dokl. 25, 2(1980), pp. 354-356
- [5] Aleksandrov V.M., Kovalenko E.V. On the contact interaction of bodies with coatings and abrasion. Sov. Phys. Dokl. 29, 4 (1984), pp. 340-342.
- [6] Aleksandrov V.M., Kovalenko E.V. Problems of Continuum Mechanics with Mixed Boundary Conditions. (in Russian). Moscow. Nauka, 1986.

- [7] Kalitvin A.S. On some class of partial integral equations in aerodynamics. (in Russian). Sost. Persp. Razv. Nauch. – Tekhn. Pod. Lipetsk. Obl(Lipetsk), 1994 pp. 210-212.
- [8] Goursat E. Cours d'Analyse Mathematique. Paris, Gautheir-Villars, 1943.
- [9] Muntz C.H. Zum dynamischen Warmeleitungs problem. Math. Z. 38, 1934, pp. 323-337.
- [10] Eshkabilov Yu.Kh. On adiscrete "three-particle" Schrodinger operator in the Hubbard model. Theor. Math. Phys., 149 (2), 2006, pp. 1497–1511.
- [11] Albeverio S., Lakaev S.N., Muminov Z. I. On the number of eigenvalues of a a model operator associated to a system of three-particles on lattices. Russ. J. Math. Phys. 2007, 14 (4), pp. 377-387.
- [12] Rasulov T.Kh. Asymptotics of the discrete spectrum of a model operator associated with a system of three particles on a lattice. Theor. and Math.Phys., 2010, 163 (1), pp.429-437.
- [13] Eshkabilov Yu.Kh., Kucharov R.R. Essential and discrete spectra of the three-particle Schrodinger operator on a lattice. Theor. Math. Phys., 170 (3), 2012,pp.341-353.
- [14] Eshkabilov Yu.Kh. Efimov effect for a 3-particle model discrete Schrodinger operator. Theor. Math. Phys., 164 (1), 2010, pp.896-90.
- [15] Appell J., Frolova E.V., Kalitvin A.S., Zabrejko P.P. Partial integral operators on  $C([a,b] \times [c,d])$ . Integral Equ. Oper. theory, 27, 1997, pp.125-140.
- [16] Likhtarnikov L.M., Vitova L.Z. On the spectrum of an integral operator with partial integrals. (in Russian), Litov. Mat. Sbornik, 15, 2(1975), pp.41-47.
- [17] Kalitvin A.S. On the spectrum of a linear operators with partial integrals and positive kernels. (in Russian), Pribl. Funk. Spektr. Theor., Leningrad, 1988, pp.43-50.
- [18] Kalitvin A.S., Zabrejko P.P. On the theory of partial integral operators. J. Integral Equations Appl., 1991,3, Num.3, pp. 351-382.
- [19] Eshkabilov Yu.Kh. Spectra of partial integral operators with a kernel of three variables. Central European J. Math., 2008, 6, Num.1, pp. 149-157.
- [20] Kalitvin A.S., Frolova E.V. Linejnye uravneniya s chastnymi integralami, Lipeck, LGPU, 2004
- [21] Kalitvin A.S. On a class of integral equations in the spase of continuous functions, -Differential Equations 42, (2006), pp. 1262-1268.
- [22] Kalitvin A.S., Kalitvin V.A. Integral'nye uravneniya Vol'terra i Vol'terra-Fredgol'ma s chastnymi integralami, Lipeck: LGPU, 2006.
- [23] Zabrejko P.P., Kalitvin A.S., Frolova E.V.: On partial integral equations in the space of continuous functions, -Differential Equations 38,4 (2002),pp.-546.
- [24] Appell J., Kalitvin A.S., Zabrejko P.P. Partial Integral operators and Integro-differential Equations. New York, Basel, 2000.
- [25] Reed M., Simon B. Methods of Modern Mathematical Physics, Vol.1, Functional Analysis. Acad. Press, New York (1972).

Department of mechanics and mathematics. National University of Uzbekistan, Vuzgorodok, 100174, Tashkent, Uzbekistan.

E-mail address: yusup62@mail.ru

Department of energetics, Tashkent State Technical University, Vuzgorodok, 100174, Tashkent, Uzbekistan.

E-mail address: arzikulov79@mail.ru

Department of mechanics and mathematics. National University of Uzbekistan, Vuzgorodok, 100174, Tashkent, Uzbekistan.

E-mail address: haydarov\_imc@mail.ru